

Capacity Bounds for the 3-dimensional (0, 1) Runlength Limited Channel

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Abstract. The capacity $C_{0,1}^{(3)}$ of a 3-dimensional (0, 1) runlength constrained channel is shown to satisfy $0.522501741838 \leq C_{0,1}^{(3)} \leq 0.526880847825$.

1 Introduction

A binary sequence satisfies a 1-dimensional (d, k) runlength constraint if there are at most k zeros in a row, and between every two consecutive ones there are at least d zeros. An n -dimensional binary array is said to satisfy a (d, k) runlength constraint, if it satisfies the 1-dimensional (d, k) runlength constraint along every direction parallel to a coordinate axis. Such an array is called *valid*. The number of valid n -dimensional arrays of size $m_1 \times m_2 \times \dots \times m_n$ is denoted by $N_{m_1, m_2, \dots, m_n}^{(d, k)}$ and the corresponding *capacity* is defined as

$$C_{d,k}^{(n)} = \lim_{m_1, m_2, \dots, m_n \rightarrow \infty} \frac{\log_2 N_{m_1, m_2, \dots, m_n}^{(d, k)}}{m_1 m_2 \dots m_n}.$$

By exchanging the roles of 0 and 1 it can be seen that $C_{0,1}^{(n)} = C_{1,\infty}^{(n)}$ for all $n \geq 1$. A simple proof of the existence of the 2-dimensional (d, k) capacities can be found in [1], and the proof can be generalized to n -dimensions.

It is known (e.g. see [2]) that the 1-dimensional (0, 1)-constrained capacity is the logarithm of the golden ratio, i.e.

$$C_{0,1}^{(1)} = \log_2 \frac{1 + \sqrt{5}}{2} = 0.694242 \dots$$

and in [3] very close upper and lower bounds were given for the 2-dimensional (0, 1)-constrained capacity. The bounds in [3] were calculated with greater precision in [4] and are further slightly improved here by us (see Remark section at end for more details), now agreeing in 9 decimal positions:

$$0.587891161775 \leq C_{0,1}^{(2)} \leq 0.587891161868 \quad (1)$$

A lower bound of $C_{0,1}^{(2)} \geq 0.5831$ was obtained in [5] by using an implementable encoding procedure known as "bit-stuffing". The known bounds on $C_{0,1}^{(2)}$ have played a useful

role in [1] for obtaining bounds on other (d, k) -constraints in two dimensions. The 3-dimensional $(0, 1)$ -constrained bounds given in the present paper can play a similar role for obtaining different 3-dimensional bounds, and are also of theoretical interest. In fact, a recent tutorial paper [6] discusses an interesting connection between run length constrained capacities in more than one dimension and crossword puzzles (based on work of Shannon from 1948). In the present paper we consider the 3-dimensional $(0, 1)$ constraint, and by extending ideas from [3] our main result is to derive (in Sections 2 and 3) the following bounds on the 3-dimensional $(0, 1)$ capacity.

Theorem 1

$$0.522501741838 \leq C_{0,1}^{(3)} \leq 0.526880847825$$

It is assumed henceforth in this paper that $d = 0$ and $k = 1$. Two valid $m_1 \times m_2$ rectangles can be put next to each other in 3 dimensions without violating the 3-dimensional $(0, 1)$ constraint if they have no two zeros in the same positions. Define a transfer matrix T_{m_1, m_2} to be an $N_{m_1, m_2}^{(0,1)} \times N_{m_1, m_2}^{(0,1)}$ binary matrix, such that the rows and columns are indexed by the valid 2-dimensional $m_1 \times m_2$ patterns, and an entry of T_{m_1, m_2} is 1 if and only if the corresponding two rectangles can be placed next to each other in 3 dimensions without violating the $(0, 1)$ constraint. Then,

$$N_{m_1, m_2, m_3}^{(0,1)} = \mathbf{1}' \cdot T_{m_1, m_2}^{m_3-1} \mathbf{1} = \mathbf{1}' \cdot T_{m_1, m_3}^{m_2-1} \mathbf{1} = \mathbf{1}' \cdot T_{m_2, m_3}^{m_1-1} \mathbf{1}$$

where $\mathbf{1}$ is the all ones column vector and prime denotes transpose. The matrix T_{m_1, m_2} meets the conditions of the Perron-Frobenius theorem [7], since it has nonnegative real elements and is irreducible (since the all one's rectangle can be placed next to any valid rectangle without violating the $(0, 1)$ constraint). Therefore the largest magnitude eigenvalue Λ_{m_1, m_2} of T_{m_1, m_2} is positive, real, and has multiplicity one. This implies that

$$\lim_{m_3 \rightarrow \infty} (N_{m_1, m_2, m_3}^{(0,1)})^{1/m_3} = \Lambda_{m_1, m_2},$$

and

$$\begin{aligned} C_{0,1}^{(3)} &= \lim_{m_1, m_2, m_3 \rightarrow \infty} \frac{\log_2 N_{m_1, m_2, m_3}^{(0,1)}}{m_1 m_2 m_3} \\ &= \lim_{m_1, m_2 \rightarrow \infty} \frac{\log_2 \lim_{m_3 \rightarrow \infty} (N_{m_1, m_2, m_3}^{(0,1)})^{1/m_3}}{m_1 m_2} \\ &= \lim_{m_1, m_2 \rightarrow \infty} \frac{\log_2 \Lambda_{m_1, m_2}}{m_1 m_2} \\ &= \lim_{m_1 \rightarrow \infty} \frac{\log_2 \lim_{m_2 \rightarrow \infty} \Lambda_{m_1, m_2}^{1/m_2}}{m_1} \\ &= \lim_{m_1 \rightarrow \infty} \frac{\log_2 \Lambda_{m_1}}{m_1}, \end{aligned} \tag{2}$$

where $\Lambda_{m_1} = \lim_{m_2 \rightarrow \infty} \Lambda_{m_1, m_2}^{1/m_2}$. The quantities $\frac{\log_2 \Lambda_{m_1, m_2}}{m_1 m_2}$ and $\frac{\log_2 \Lambda_{m_1}}{m_1}$ can be viewed as capacities corresponding to 3-dimensional arrays with two fixed sides (lengths m_1 and m_2), and one fixed side (length m_1), respectively.

Upper and lower bounds on the 3-dimensional capacity can be computed directly from the inequalities (similar to the 2-dimensional case, as noted in [4])

$$\frac{\log_2 \Lambda_{m_1, m_2}}{(m_1 + 1)(m_2 + 1)} \leq C_{0,1}^{(3)} \leq \frac{\log_2 \Lambda_{m_1, m_2}}{m_1 m_2}$$

but these do not yield particularly tight bounds for values of m_1 and m_2 that result in reasonable space and time complexities (e.g. Table 1 shows that the eigenvalues Λ_{m_1, m_2} correspond to matrices with more than 40 million elements when roughly $m_1 m_2 \geq 20$). The upper and lower capacity bounds derived in this paper agree to within ± 0.002 and were computed using less than 100 Mbytes of computer memory.

2 Lower bound on $C_{0,1}^{(3)}$

To derive a lower bound on $C_{0,1}^{(3)}$ we generalize a method of Calkin and Wilf [3]. Since T_{m_1, m_2} is a symmetric matrix, the Courant-Fischer Minimax Theorem [8, pg. 394] implies that

$$\Lambda_{m_1, m_2}^p \geq \frac{\mathbf{x}' \cdot T_{m_1, m_2}^p \mathbf{x}}{\mathbf{x}' \cdot \mathbf{x}} \quad (3)$$

for any nonzero vector \mathbf{x} and any integer $p \geq 0$. Choosing $\mathbf{x} = T_{m_1, m_2}^q \mathbf{1}$ for any integer $q \geq 0$ gives

$$\Lambda_{m_1, m_2}^p \geq \frac{\mathbf{1}' \cdot T_{m_1, m_2}^{p+2q} \mathbf{1}}{\mathbf{1}' \cdot T_{m_1, m_2}^{2q} \mathbf{1}} = \frac{\mathbf{1}' \cdot T_{m_1, p+2q+1}^{m_2-1} \mathbf{1}}{\mathbf{1}' \cdot T_{m_1, 2q+1}^{m_2-1} \mathbf{1}} \quad (4)$$

Thus,

$$\begin{aligned} 2^{pC_{0,1}^{(3)}} &= \left(\lim_{m_1, m_2 \rightarrow \infty} \Lambda_{m_1, m_2}^{1/(m_1 m_2)} \right)^p = \lim_{m_1 \rightarrow \infty} \left(\lim_{m_2 \rightarrow \infty} \Lambda_{m_1, m_2}^{p/m_2} \right)^{1/m_1} \\ &\geq \lim_{m_1 \rightarrow \infty} \left(\frac{\Lambda_{m_1, p+2q+1}}{\Lambda_{m_1, 2q+1}} \right)^{1/m_1} = \frac{\lim_{m_1 \rightarrow \infty} \Lambda_{m_1, p+2q+1}^{1/m_1}}{\lim_{m_1 \rightarrow \infty} \Lambda_{m_1, 2q+1}^{1/m_1}} = \frac{\Lambda_{p+2q+1}}{\Lambda_{2q+1}} \quad (5) \end{aligned}$$

and therefore for any odd integer $r \geq 1$ and any integer $z > r$,

$$C_{0,1}^{(3)} \geq \frac{1}{z-r} \log_2 \left(\frac{\Lambda_z}{\Lambda_r} \right) \quad (6)$$

This lower bound on $C_{0,1}^{(3)}$ is analogous to a 2-dimensional bound in [3], but Λ_z and Λ_r are not eigenvalues associated with transfer matrices of 2-dimensional arrays here, and cannot easily be computed as in the 2-dimensional case. Instead, we obtain a lower bound on Λ_z and an upper bound on Λ_r . From (4) and (5) a lower bound on Λ_z is

$$\Lambda_z = \lim_{m_2 \rightarrow \infty} \Lambda_{z, m_2}^{1/m_2} \geq \lim_{m_2 \rightarrow \infty} \left(\frac{\mathbf{1}' \cdot T_{z, v}^{m_2-1} \mathbf{1}}{\mathbf{1}' \cdot T_{z, u}^{m_2-1} \mathbf{1}} \right)^{1/((v-u)m_2)} = \left(\frac{\Lambda_{z, v}}{\Lambda_{z, u}} \right)^{1/(v-u)}$$

where u is an arbitrary positive odd integer, $v > u$, and $\Lambda_{z,v}$ and $\Lambda_{z,u}$ are the largest eigenvalues of the transfer matrices $T_{z,v}$ and $T_{z,u}$, respectively.

To find an upper bound on the quantity Λ_r for a given r , we apply a modified version of a method in [3]. We say that a binary matrix satisfies the $(0, 1)$ cylindrical constraint if it satisfies the usual 2-dimensional $(0, 1)$ constraint after joining its leftmost column to its rightmost column (i.e. the left and right columns can be put next to each other without violating the $(0, 1)$ constraint). A binary matrix satisfies the $(0, 1)$ toroidal constraint if it satisfies the usual 2-dimensional $(0, 1)$ constraint after both joining its leftmost column to its rightmost column, and its top row to its bottom row.

Proposition 1 Let s be a positive even integer and let T_{m_1, m_2} be the transfer matrix whose rows and columns are indexed by all $(0, 1)$ -constrained $m_1 \times m_2$ rectangles. Let $B_{m_1, s}$ denote the transfer matrix whose rows and columns are indexed by all cylindrically $(0, 1)$ -constrained $m_1 \times s$ rectangles. Then,

$$\text{Trace}[T_{m_1, m_2}^s] = \mathbf{1}' \cdot B_{m_1, s}^{m_2 - 1} \mathbf{1}.$$

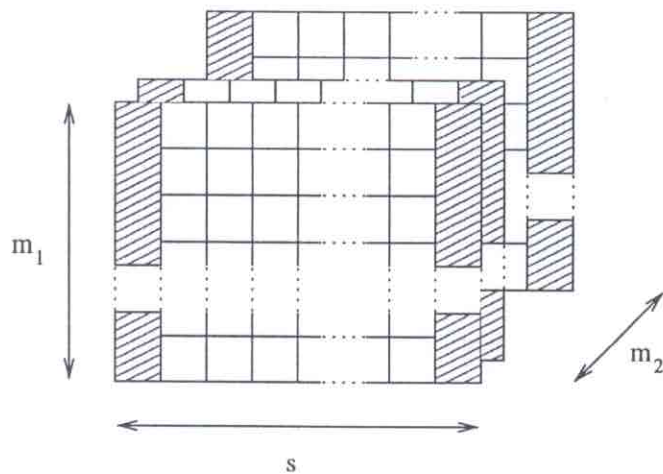


Fig. 1. Cylindrically $(0, 1)$ -constrained $m_1 \times s$ rectangles used to build cylindric $m_1 \times m_2 \times s$ arrays

For every positive integer m_1 and m_2 , and every even positive integer s , the matrix T_{m_1, m_2}^s has nonnegative eigenvalues and thus any one of its eigenvalues is upper bounded by its trace. Hence,

$$\Lambda_{m_1, m_2} \leq \text{Trace} [T_{m_1, m_2}^s]^{1/s} = (\mathbf{1}' \cdot B_{m_1, s}^{m_2 - 1} \mathbf{1})^{1/s} \quad (7)$$

which gives the following upper bound on Λ_r :

$$\Lambda_r = \lim_{m_2 \rightarrow \infty} \Lambda_{r, m_2}^{1/m_2} \leq \lim_{m_2 \rightarrow \infty} (\mathbf{1}' \cdot B_{r, s}^{m_2 - 1} \mathbf{1})^{1/m_2} = \xi_{r, s}^{1/s}, \quad (8)$$

where $\xi_{r,s}$ is the largest eigenvalue of $B_{r,s}$ (note that $B_{r,s}$ satisfies the Perron-Frobenius theorem for the same reasons as for T_{m_1,m_2} in Section 1).

The lower bound on $C_{0,1}^{(3)}$ in (6) can now be written as

$$C_{0,1}^{(3)} \geq \frac{1}{z-r} \log_2 \left(\frac{\left(\frac{A_{z,v}}{A_{z,u}} \right)^{1/(v-u)}}{\xi_{r,s}^{1/s}} \right) \quad \begin{array}{l} r \text{ and } u \text{ odd, } s \text{ even} \\ z > r \geq 1 \\ v > u \geq 1 \\ s \geq 2 \end{array} \quad (9)$$

To obtain the best possible lower bound, the right hand side of (9) should be maximized over all acceptable choices of r , z , u , v , and s , subject to the numerical computability of the quantities $A_{z,v}$, $A_{z,u}$, and $\xi_{r,s}$. Table 1 shows the largest eigenvalues of various transfer matrices which were numerically computable. From this table, the best parameters we could find for the lower bound in (9) on the capacity were $r = 3$, $z = 4$, $u = 5$, $v = 6$, and $s = 10$, yielding

$$C_{0,1}^{(3)} \geq \frac{1}{4-3} \log_2 \frac{\frac{9346.35893701}{2102.73425568}}{(80481.0598379)^{1/10}} \geq 0.522501741838.$$

3 Upper bound on $C_{0,1}^{(3)}$

Proposition 2 Let s_1 and s_2 be positive even integers and let B_{s_1,s_2}^* denote the transfer matrix whose rows and columns are indexed by all toroidally $(0,1)$ -constrained $s_1 \times s_2$ rectangles. If ξ_{s_1,s_2}^* is the largest eigenvalue of B_{s_1,s_2}^* , then $C_{0,1}^{(3)} \leq \frac{1}{s_1 s_2} \log_2 \xi_{s_1,s_2}^*$.

Note that $B_{2,s_2} = B_{2,s_2}^*$ and thus $\xi_{2,s_2} = \xi_{2,s_2}^*$. The best parameters we were able to find (from Table 1) were $s_1 = 4$ and $s_2 = 6$, and the resulting eigenvalue gave the following upper bound:

$$C_{0,1}^{(3)} \leq \frac{1}{24} \log_2 6405.69924332 \leq 0.526880847825.$$

4 Remark

Direct computation of eigenvalues using standard linear algebra algorithms generally requires the storage of an entire matrix. This severely restricts the matrix sizes allowable, due to memory constraints on computers. By exploiting the fact that our matrices are all binary, symmetric, and easily computable, we were able to obtain the largest eigenvalues of much larger matrices. Specifically, the eigenvalues used to obtain the capacity bounds in Theorem 1 were computed using the "power method" [8, pg. 406]. Similarly, we obtained the upper bound in (1) with the power method (computing $A_{1,21}$, $A_{1,23}$, and $\xi_{1,24}$). Originally these bounds were computed in [3] as $0.587891161 \leq C_{0,1}^{(2)} \leq 0.588339078$ (computing $A_{1,13}$, $A_{1,15}$, and $\xi_{1,6}$) and were later improved in [4] (computing $A_{1,13}$, $A_{1,14}$, and $\xi_{1,14}$) to $0.587891161775 \leq C_{0,1}^{(2)} \leq 0.587891494943$. The lower bound in (1) is from [4].

a	b	$\Lambda_{a,b}$	rows of $T_{a,b}$	$\xi_{a,b}$	rows of $B_{a,b}$	$\xi_{a,b}^*$	rows of $B_{a,b}^*$	
1	1	1.61803398875	2					
	2	2.41421356237	3	2.41421356237	3			
	3	3.63138126040	5					
	4	5.45770539597	8	5.15632517466	7			
	5	8.20325919376	13					
	6	12.3298822153	21	11.5517095660	18			
	7	18.5324073775	34					
	8	27.8550990963	55	26.0579860919	47			
	9	41.8675533183	89					
	10	62.9289457252	144	58.8519350815	123			
	11	94.5852312050	233					
	12	142.166150393	377	132.947794048	322			
	13	213.682559741	610					
	14	321.175161677	987	300.345852027	843			
	15	482.741710897	1597					
	16	725.584002895	2584	678.525669346	2207			
	17	1090.58764423	4181					
	18	1639.20566742	6765	1532.89283597	5778			
	19	2463.80493521	10946					
	20	3703.21728345	17711	3463.03987027	15127			
	21	5566.11363689	28657					
	22	8366.13642876	46368	7823.53857819	39603			
	23	12574.7053170	75025					
	24	18900.3867144	121393	17674.5747630	103682			
2	2	5.15632517466	7	5.15632517466	7	5.15632517466	7	
	3	11.1103016575	17					
	4	23.9250625386	41	21.9287654025	35	21.9287654025	35	
	5	51.5229210280	99					
	6	110.954925971	239	100.236549238	199	100.236549239	199	
	7	238.942175857	577					
	8	514.563569622	1393	463.203410887	1155	463.203410887	1155	
	9	1108.11608218	3363					
	10	2386.33538059	8119	2146.04060032	6727	2146.04060032	6727	
	11	5138.98917320	19601					
	12	11066.8474924	47312	9949.63685703	39203	9949.63685703	39203	
	3	3	34.4037405361	63				
4		106.439377528	227	94.2548937790	181			
5		329.331697608	827					
6		1018.97101980	2999	884.498791440	2309			
7		3152.75734322	10897					
8		9754.81971205	39561	8421.60680806	30277			
9		30181.9963196	143677					
10		93384.9044989	521721	80481.0598378	398857			
4		4	473.069084944	1234	404.943621498	933	355.525781764	743
		5	2102.73425567	6743				
	6	9346.35893702	36787	7799.87080772	26660	6405.69924332	18995	

Table 1. Largest eigenvalues of $T_{a,b}$, $B_{a,b}$, and $B_{a,b}^*$ are $\Lambda_{a,b}$, $\xi_{a,b}$, and $\xi_{a,b}^*$.

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